

## Short Communication

# Theoretical Analysis of the Optical Properties of Gold Nanoparticles Using DDA Approximation

H. Akherat Doost, M. H. Majles Ara\* and E. Koshki

Department of Physics, Kharazmi University, Mofateh Ave, Tehran, Iran.

(\*) Corresponding author: majlesara@gmail.com

(Received: 09 November 2015 and Accepted: 07 August 2016)

### **Abstract**

*This article describes a study, using numerical simulation, of the optical properties of nano particles as a function of their size. Many methods introduced to simulate and calculate the interaction of light and particle, such as Mie analysis, boundary element and finite element methods. The Discrete Dipole Approximation (DDA), wherein a target geometry is modeled as a finite array of polarizable elements and DDSCAT simulation, are employed for determining extinction, absorption and scattering cross sections by gold nano particles with 15, 35, 55, 75, 100, 125 and 150 nanometers in diameter in different wavelengths. In this method, the particle is represented as a cubic array of electric dipoles. The results show that Plasmon resonance dependent on nano particles size. In other project, the angular distribution of different sizes of gold nano particles are calculated and scattering of these particles are compared. This study can be helpful in designing appropriately shaped nano particles for in-vivo applications like photo-thermal cancer treatment and optical sensors.*

**Keywords:** Absorption and scattering of gold nanoparticles, Discrete Dipole Approximation (DDA), DDSCAT simulation.

## 1. INTRODUCTION

Geometry of nano particles and efficient modeling of light scattered from small features and particles on surfaces are primary factor and important facility for detecting and analyzing how incident light is scattered and absorbed by particles. Recently physical properties of different sizes of metallic nanoparticles and optical models of nano- metrology systems based on laser are theoretically analyzed [1, 2]. Prediction of evanescent wave scattering by particles on a surface is a challenging problem and knowledge of nanoparticle-light interactions and their cross sections calculation are very helpful in many studies such as soot agglomerates in combustion systems, interstellar dust grains in astrophysics, biological cells in biomedical applications and aerosols in

environmental research and atmospheric science. Development and future of more sensitive inspection systems will require a flexible and reliable scattering model. One of the greatest triumphs of classical physics was achieved in 1908, when Gustav Mie presented a solution to Maxwell's equations that described the extinction spectra (extinction=scattering + absorption) of spherical particles of arbitrary size. The T-matrix method combined with the normal incidence approximation (NIA) has been used to predict scattering by a sphere on a surface illuminated by a propagating [19, 20] and an evanescent [8, 18, 21, 23] wave. An exact solution for the problem of multiple spheres on a surface was also proposed by Mackowski [11]. For characterization

purposes, the forward model should be able to accommodate arbitrarily shaped objects on a surface. Numerical approaches such as the null-field method with discrete sources and the finite-difference time-domain method [12] could be employed for that purpose. We have used the discrete-dipole approximation (DDA) method to compute all cross sections of nanoparticles. Discrete dipole approximation (DDA) is a well-established and widely accepted method of solving light scattering by particles from direct propagating illumination. DDA was proposed by Purcell and Pennypacker (PP) [14], who replaced the scatterer by a set of targets, these dipoles interact with each other and the incident field and each of which attains a dipole moment due to both the incident beam and the other oscillating dipoles within the target. The DDA is a popular method in the light-scattering community and it has been reviewed by several authors. The DDA was further developed by Draine and coworkers [3, 4, 5, 6, 7], who describes the current state of the DDA and its historical development. It also explains the equivalence of the DDA and methods based on the volume integral equation formulation and popularized the method by developing a publicly available computer code DDSCAT [6]. Later it was shown that the DDA also can be derived from the integral equation for the electric field. This derivation was apparently first performed by Goedecke and O'Brien [9] and further developed by others (see, for instance, [10, 13, 15, 24] and certain DDA applications are concerned with very prolate or oblate particles [16, 17, 22]). In this paper, at first, the extinction, scattering and absorption cross sections for different sizes of gold nanoparticles were determined using DDA technique and then a comparison of plasmon resonance peaks for all nanoparticles was made.

## 2. DIPOLE DISCRETE APPROXIMATION METHOD

In this method, the particle is represented as a cubic array of electric dipoles. Assume the number of dipoles is  $N$ . The  $j$ th dipole has a polarizability  $\alpha_j$ , and the total electric field at the dipole is  $E_j$ . Therefore, the polarization  $P_j$  in response to the total electromagnetic field can be described by

$$P_j = \alpha_j E_j = \left( E_{inc,j} - \sum_{k \neq j} A_{jk} \cdot P_k \right) \quad (1)$$

where  $E_j$  is the local electric field, which is the sum of the incident field and the retarded radiation fields from the other dipoles, The incident field is assumed to be a plane wave and is in the form of  $E_{inc,j} = E_0 e^{i\mathbf{k} \cdot \mathbf{r}_j}$  where  $\mathbf{k}$  is the wave vector.  $\mathbf{A}$  is the dipole-dipole interaction matrix,

$$A_{jk} P_k = \frac{\exp(i\mathbf{k} \cdot \mathbf{r}_{jk})}{r_{jk}^3} \left\{ k^2 \overline{r_{jk}} \times (\overline{r_{jk}} \times \overline{P_k}) + \frac{1 - i\mathbf{k} \cdot \overline{r_{jk}}}{r_{jk}^2} \left[ r_{jk}^2 \overline{P_k} - 3\overline{r_{jk}} (\overline{r_{jk}} \cdot \overline{P_k}) \right] \right\} \quad (2)$$

Here,  $k$  is the wave number, and  $\overline{r_{jk}} = \overline{r_j} - \overline{r_k}$ . Let  $A_{jj} = \alpha_j^{-1}$  so that we can obtain:

$$\mathbf{A} \mathbf{P} = \mathbf{E}_{inc} \quad (3)$$

$\mathbf{P}$  and  $\mathbf{E}_{inc}$  are  $3N$ -dimensional vectors, and  $\mathbf{A}$  is a  $3N \times 3N$  matrix. Equation 3 can be solved iteratively by means of the complex conjugate gradient (CCG) algorithm in combination with the fast Fourier transform technique. When the polarization of each dipole is known, the extinction, absorption and scattering cross-sections can be computed using

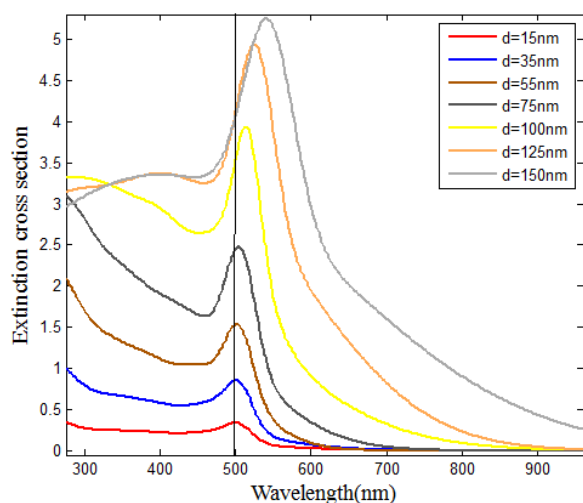
$$C_{ext} = \frac{4\pi k}{|\mathbf{E}_{inc}|^2} \sum_{k \neq j} \text{Im}(\mathbf{E}_{inc,j}^* \cdot \mathbf{P}_j) \quad (4)$$

$$C_{abs} = \frac{4\pi k}{|\mathbf{E}_0|^2} \times \sum_{j=1}^N \left\{ \text{Im} \left[ \mathbf{P}_j \cdot \alpha_j^{-1} \cdot \mathbf{P}_j \right] - \frac{2}{3} k^3 |\mathbf{P}_j|^2 \right\} \quad (5)$$

$$C_{sca} = C_{ext} - C_{abs} \quad (6)$$

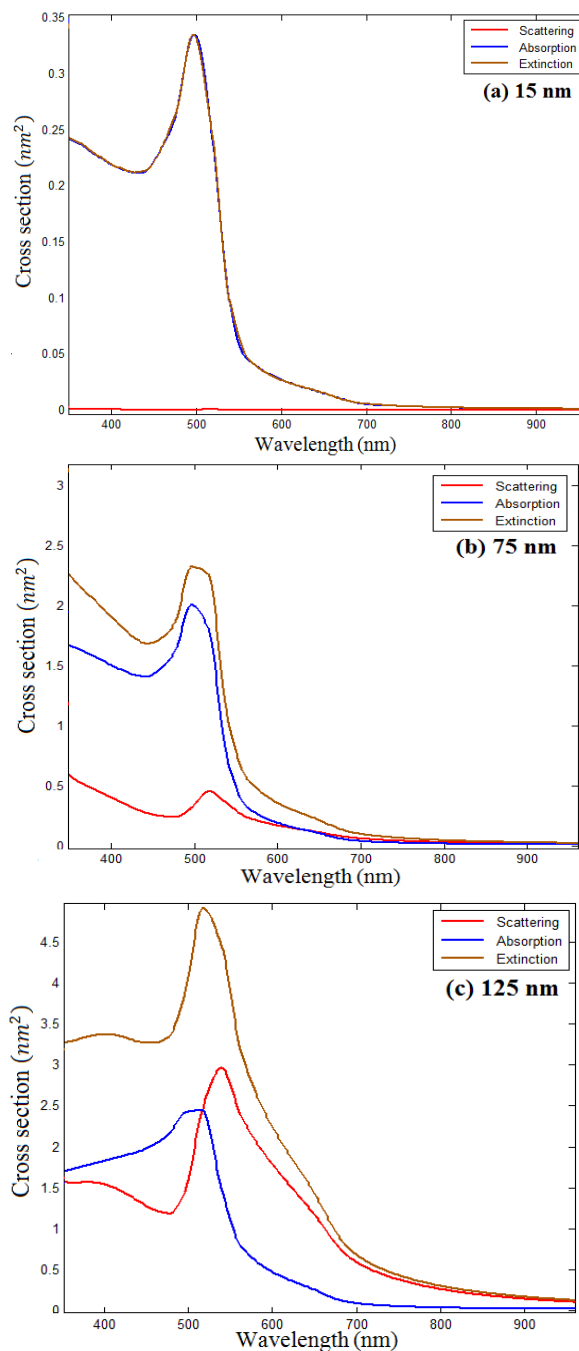
### 3. RESULTS AND DISCUSSIONS

It is possible to calculate scattering, absorption and extinction cross-sections of different sizes of gold nanoparticles by using DDA (Fig.1, 3-4). All the simulation results for particles of arbitrary geometries and different sizes reported in this work were obtained. Extinction is the sum of scattering and absorption which is a measure of the light that would not be received by the detector that is collinear with the nanoparticle and source.



**Figure 1.** Comparison extinction spectra of 15,35,55,75,100,125 and 150 nm diameters for Au nanoparticles.

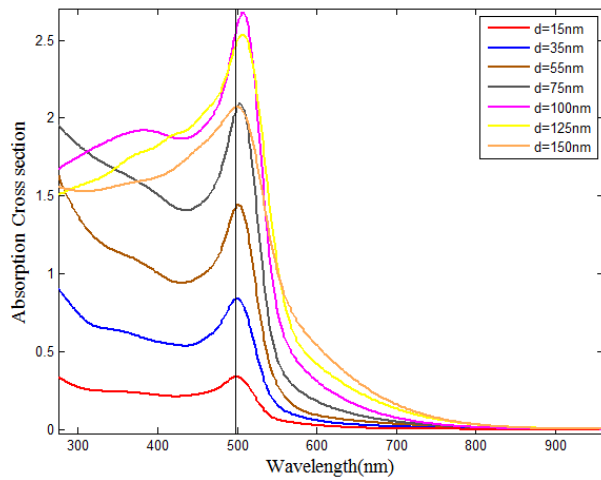
The extinction peak for Au nanoparticles occurs at the surface plasmon resonance (SPR) frequency. At this frequency, conduction electrons oscillate coherently. SPR frequencies of gold nanoparticles occur in the visible part of the spectrum, which is useful for nanoantennas, bistable devices and solar cell applications. Figure 1 shows that the extinction efficiency (i.e. extinction coefficient per unit cross sectional area of the nanoparticle) increases linearly with particle size up to a size of 150 nm. For particle sizes lower than 55 nm, the peak wavelength remains nearly constant (~520 nm) but for particle sizes higher than 55 nm, shifts to higher wavelengths. Peak wavelength values for spheres with size of 15 nm, 75 nm and 125 nm are 495 nm, 500 nm and 540 nm, respectively.



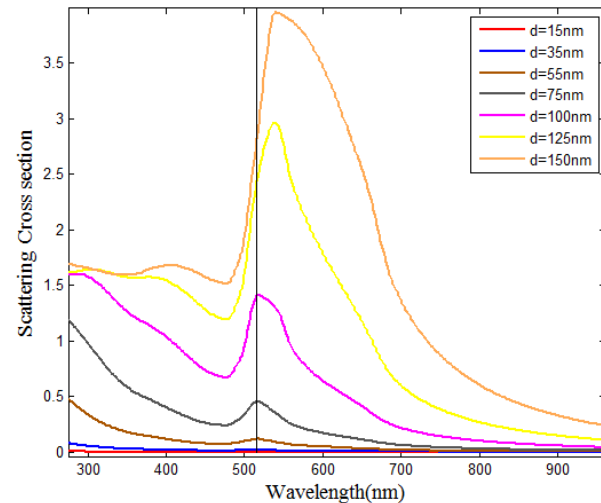
**Figure 2.** Variation of extinction, absorption and scattering efficiencies of gold nanoparticles of diameter (a) 15nm, (b) 75nm and (c) 125 nm.

Also, the computed spectra deviate from the dipolar approximation at larger sizes. These plots validate our implementation of the DDSCAT code for calculating the optical properties of nanostructures. Figures 2 (a-c) shows the variation in contributions from absorption and scattering to the overall extinction efficiency as a function of particle size for

spherical gold nanoparticles. It is seen that for smaller nanoparticles, the extinction efficiency is predominantly governed by absorption efficiency. As the size of nanoparticle increases, scattering efficiency becomes more dominant. As Well As comparison of different sizes of gold nanoparticles absorption and scattering cross sections were shown in figures 3 and 4, respectively. It is clear that noble metal nanoparticles such as gold exhibit a strong optical extinction peak in the visible range of spectrum.



**Figure 3.** Comparison absorption spectra of 15,35,55,75,100,125 and 150 nm diameters for Au nanoparticles.



**Figure 4.** Comparison scattering spectra of 15,35,55,75,100,125 and 150 nm diameters for Au nanoparticles.

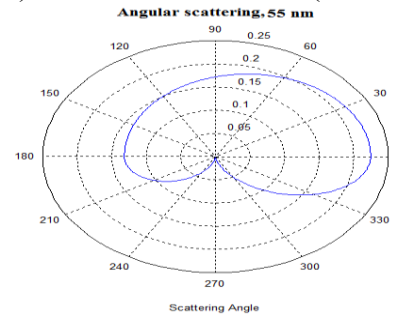
The scattered far field in spherical coordinates ( $E_{s\theta}, E_{s\phi}$ ) for a unit amplitude incident field is given by

$$E_{s\theta} = \frac{e^{ikr}}{-ikr} \cos\phi \cdot S_2(\cos\theta)$$

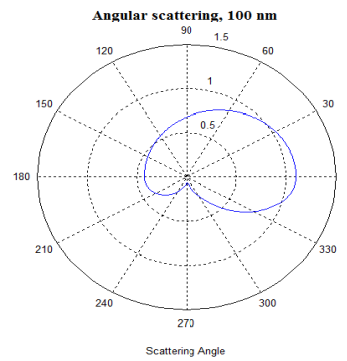
$$E_{s\phi} = \frac{e^{ikr}}{ikr} \sin\phi \cdot S_1(\cos\theta)$$

$E_{s\theta}$  is the scattered far-field component in the scattering plane, defined by the incident and scattered directions, and  $E_{s\phi}$  is the orthogonal component. The angle  $\phi$  is the angle between the incident electric field and the scattering plane.

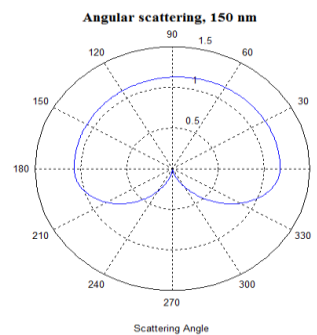
With respect to the application in a device, placing nanoparticles at an interface is a more realistic configuration. This also plays an important role when scattering efficiencies are judging. In the following, we compute scattering intensities to display results as a polar diagram of  $\theta$  in both the upper half circle ( $0 < \theta < \pi$ ) and lower half circle ( $\pi < \theta < 2\pi$ ).



(a)



(b)



(c)

**Figure 5.** Angular scattering diagram of upper and lower half circle for 55, 100, 150 diameters of Ag nanoparticles. It can be observed that scattering in the backward hemisphere is exiguity larger than that in the forward hemisphere.

#### 4. CONCLUSION

In this paper, the optical spectra of different sizes of gold nanoparticles have been simulated using DDASCAT code. Results showed that with increasing in particle size, Coulomb's force in electrons motion was smaller and electrons oscillated with low energy. It is clear that for smaller nano spheres, only absorption efficiency contributes significantly towards extinction efficiency, but with an increase

in size the contribution from scattering dominates, which was seen in Figure 2, and a linear relationship was also seen between the extinction efficiency and particle size for the smaller nano particles. Comparison between nano particles spectra showed that with increasing in size, SPR peaks had red shift and this has extent application in photonic such as nanoantennas and optical bistability devices. Finally angular distributions of Au nano particles in 55, 100, 150 nm diameters were computed that can be helpful in designing appropriate target nano particle geometry for in-vivo applications like photo-thermal cancer treatment, optical limiting devices and optical sensors.

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